

# ***Muon reconstruction and optimal event classification in AMANDA***

Gary C. Hill

University of Wisconsin- Madison

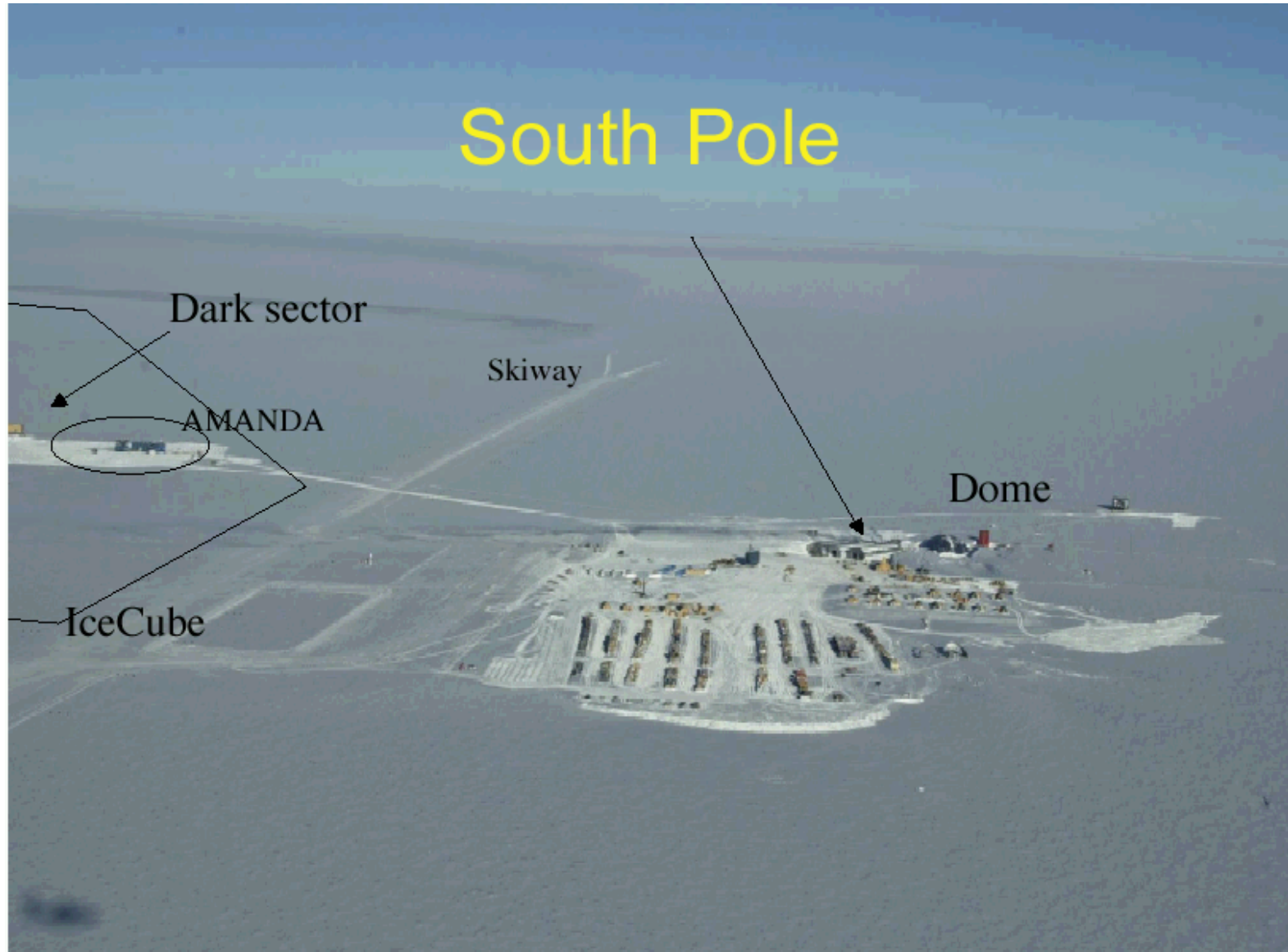
for the AMANDA Collaboration

INTERNATIONAL WORKSHOP ON  
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# Outline

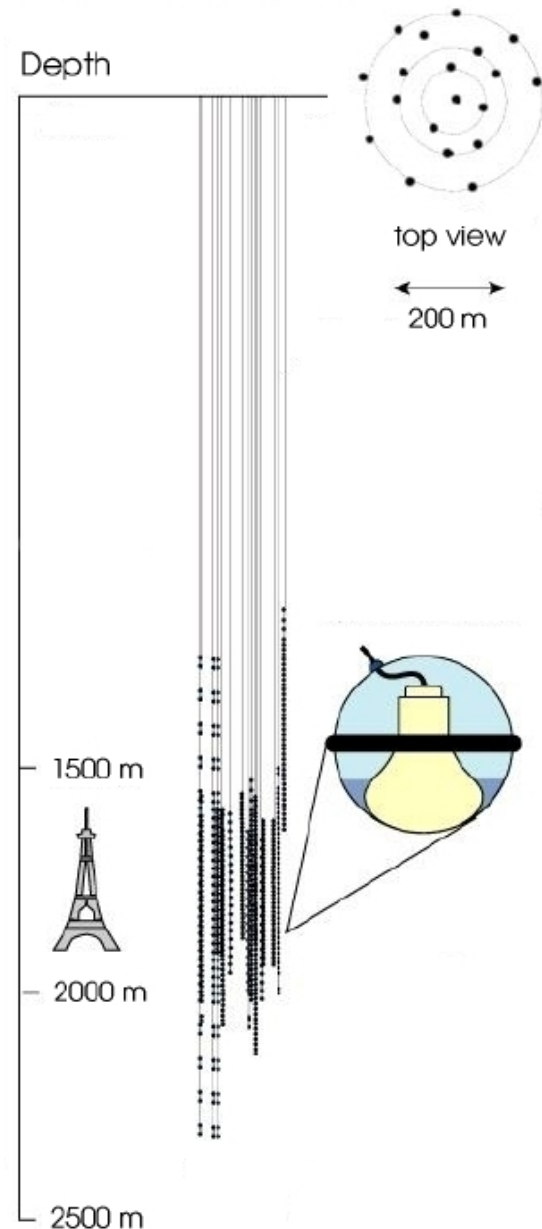
- Maximum likelihood reconstruction in the presence of light scattering in ice
- Optimal event classification - our job of reconstruction isn't done until we've assigned the most probable origin (background, signal) to an observed event
  - Zenith weighted "Bayesian" event reconstruction
  - Optimal classification with modern machine learning methods
- Try to give a unifying theme to the problem of reconstruction and event classification

# AMANDA-II Location



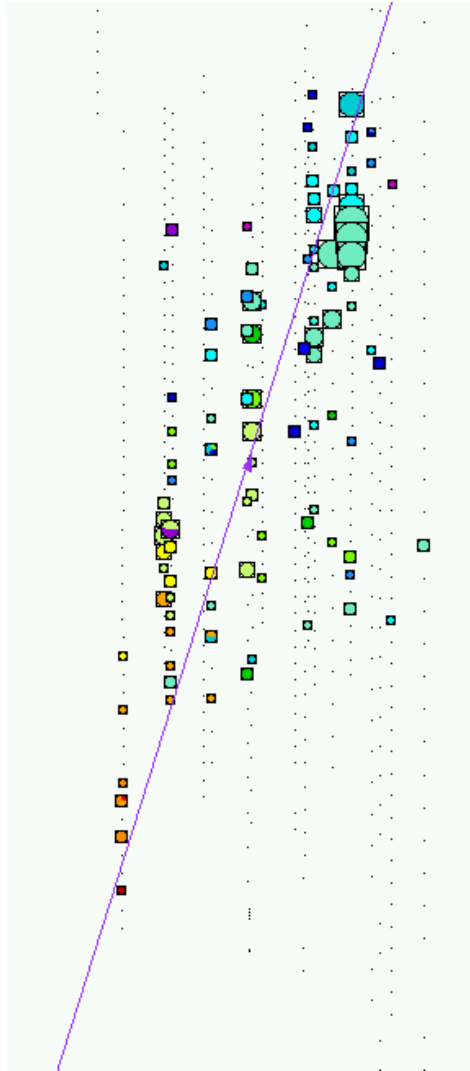
# AMANDA-II Experiment

## AMANDA-II

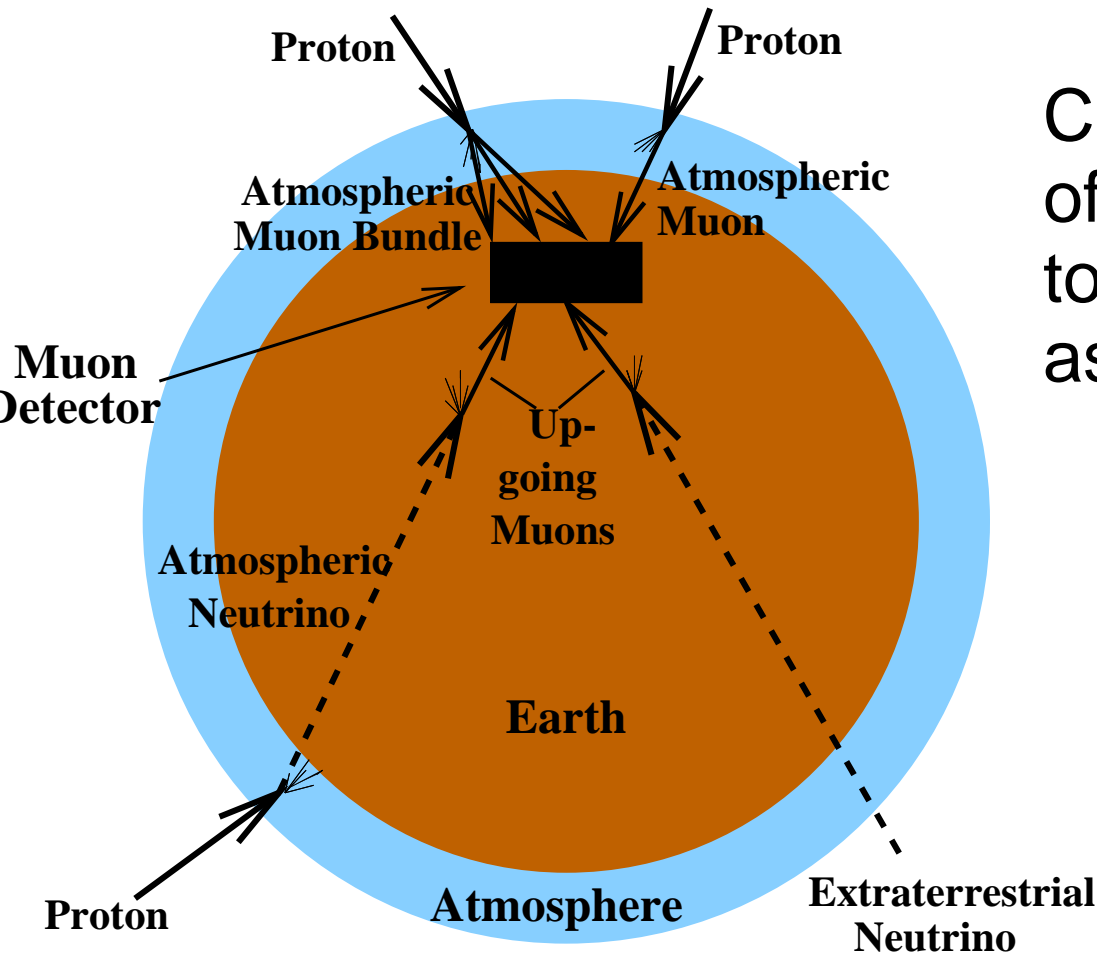


- 19 strings
- 677 Optical Modules (OM)
- 200 meters diameter
- 500 meters tall
- completed in 1999
- 1997-99 AMANDA-B10
  - 10 strings, 300 OM

# Experimental $\nu_\mu$ event



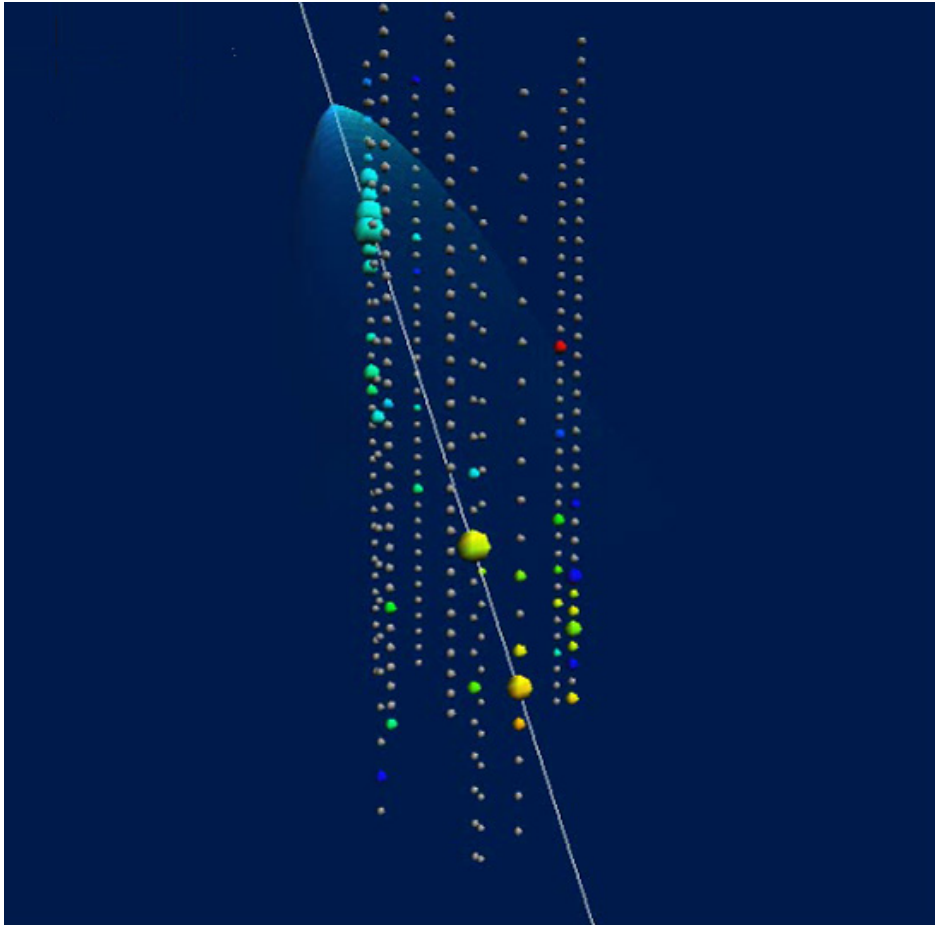
# Potential (muon) event origins?



Consider three types of hypothetical origin to which we will try to assign an event

- Downgoing muons
- Upgoing atmospheric neutrinos
- Upgoing extraterrestrial neutrinos

# *Reconstruction principle*



- Cherenkov photons are detected by PMTs
- tracks are reconstructed by maximum likelihood method of photon arrival times

# ***Muon track reconstruction***

- Cherenkov photons from the muons are recorded by the array optical modules
- each module records photon arrival times and amplitudes
- an event  $E$  is described by a vector of times and amplitudes of all the hits :

$$E \equiv \{t_1, \dots, t_n; \rho_1, \dots, \rho_n\}$$

- Wish to fit a track hypothesis :

$$H \equiv \{x, y, z, \theta, \phi\}$$

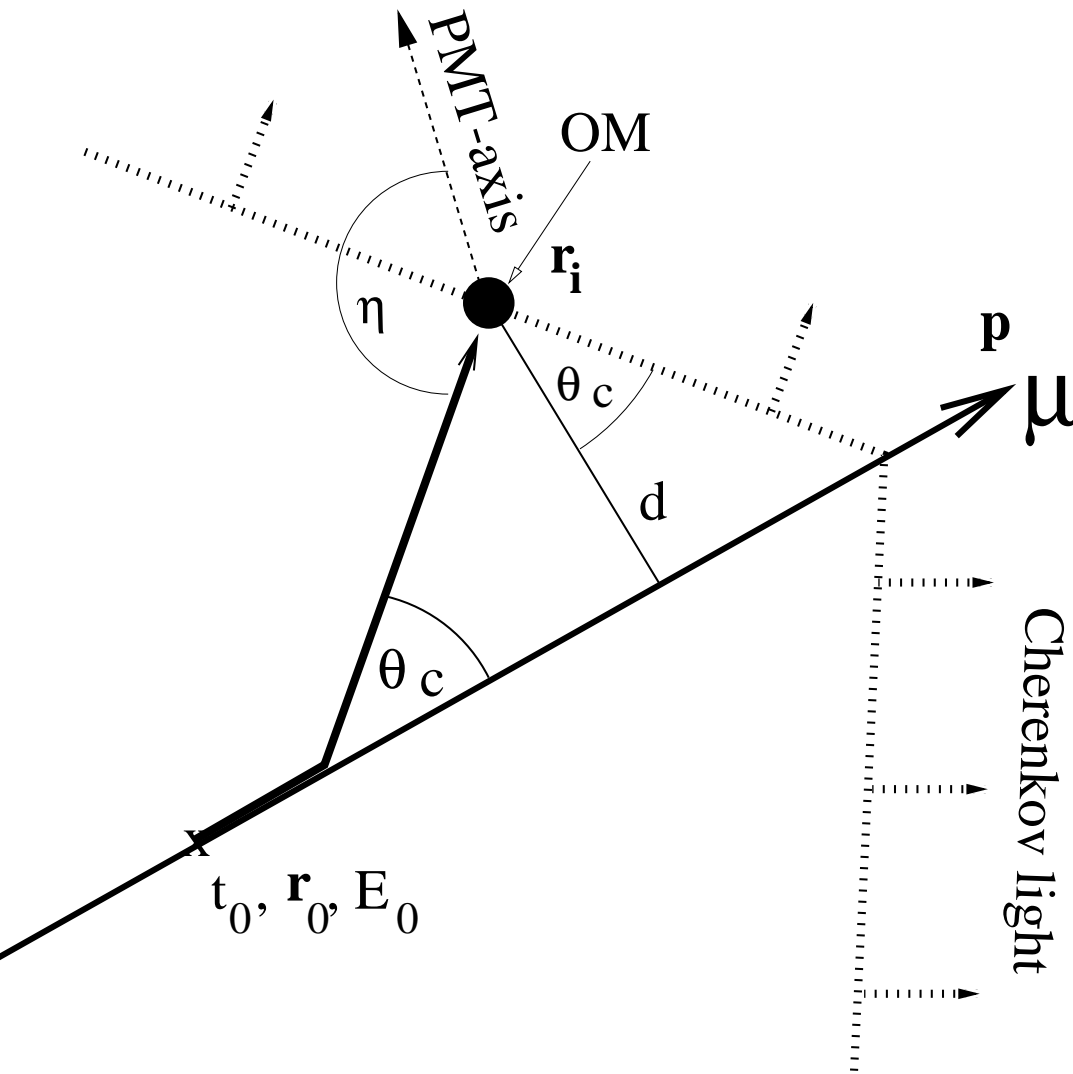


# *Fitting a muon track hypothesis to the event information*

- To connect the event  $E \equiv \{t_1, \dots, t_n; \rho_1, \dots, \rho_n\}$  and the track hypothesis  $H \equiv \{x, y, z, \theta, \phi\}$  we need the likelihood function

$$\mathcal{L}(E \equiv \{t_1, \dots, t_n; \rho_1, \dots, \rho_n\} \mid H \equiv \{x, y, z, \theta, \phi\})$$

# Muon track Cherenkov cone geometry



- Given a track hypothesis we can calculate the expected photon arrival times from an unscattered Cherenkov cone

# *Likelihood reconstruction in the absence of scattering*

- Expected photon arrival times derived from Cherenkov geometry smeared with Gaussian PMT jitter
- Straightforward form of  $p(\text{times} \mid \text{track})$

$$\mathcal{L} = \prod_{\text{OMs}} p(\text{time}_{\text{OM}} \mid \text{track})$$

- Essentially  $\chi^2$  fit
- This method insufficient in ice with scattering
- Need to use a likelihood with full photon propagation information

## *Detemining the PMT time residuals*

- Time residual is the delay in photon arrival time after the expected “direct” Cherenkov arrival time
- Full photon propagation simulation (e.g. PTD (Albrecht Karle), Photonics (Ped Miocinović)) used to tabulate residuals as a function of possible muon tracks
- These tables can be used as the reconstruction likelihood

## Analytic form - the Pandel function

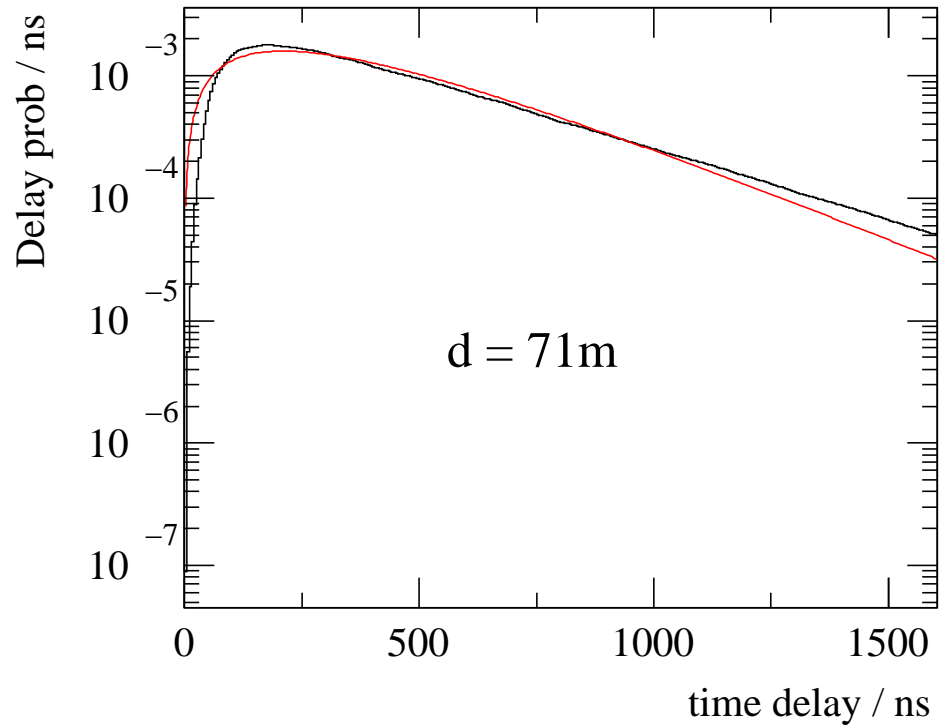
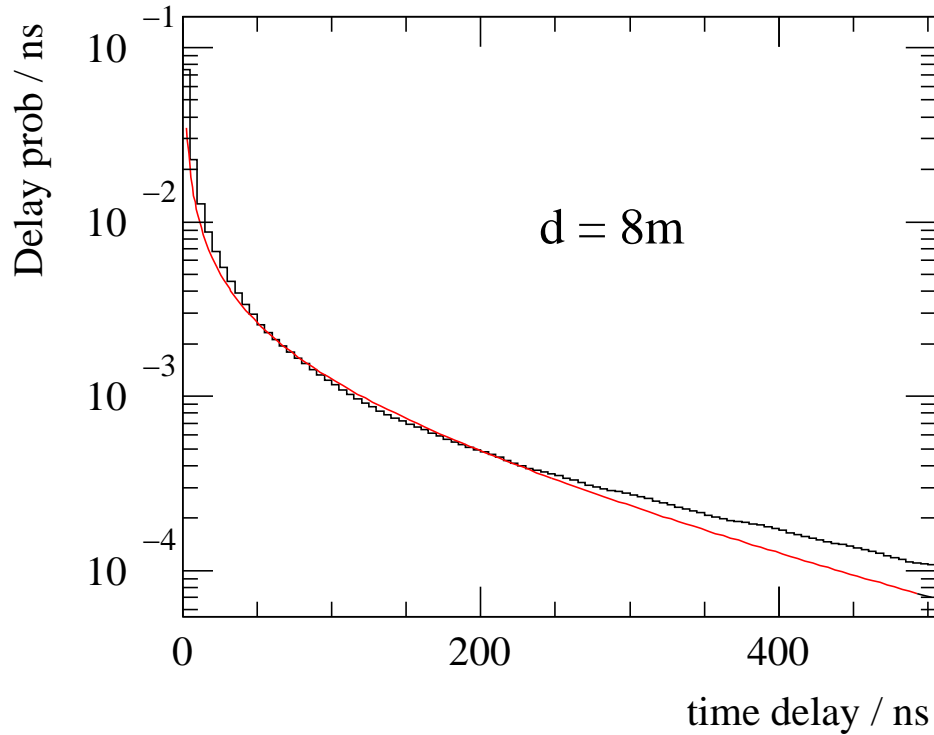
- Dirk Pandel (diploma student at DESY-Zeuthen in mid-90's) solved the propagation equations of light in the presence of absorption and scattering and found an analytic form for the time residuals

$$p(t_{\text{res}}) \equiv \frac{1}{N(d)} \frac{\tau^{-(d/\lambda)} \cdot t_{\text{res}}^{(d/\lambda-1)}}{\Gamma(d/\lambda)} \cdot e^{-\left(t_{\text{res}} \cdot \left(\frac{1}{\tau} + \frac{c_{\text{medium}}}{\lambda_a}\right) + \frac{d}{\lambda_a}\right)}$$
$$N(d) = e^{-d/\lambda_a} \cdot \left(1 + \frac{\tau \cdot c_{\text{medium}}}{\lambda_a}\right)^{-d/\lambda}$$

# *Fitting the Pandel function free parameters*

- The free parameters are fitted to make the Pandel form of the residual distributions match the full photon simulation
- Gives an analytic form that can be used in the reconstruction algorithm
- Need to add PMT jitter - old method was a simple patching of a Gaussian with the Pandel
- Recently an analytic form of the convolution of the Pandel with a Gaussian was found (George Japaridze)

# *Time residuals - full simulation and Pandel fit*



**Pandel fit (red) to photon tables (black)**

# Reconstructing an event!

- Our likelihood function

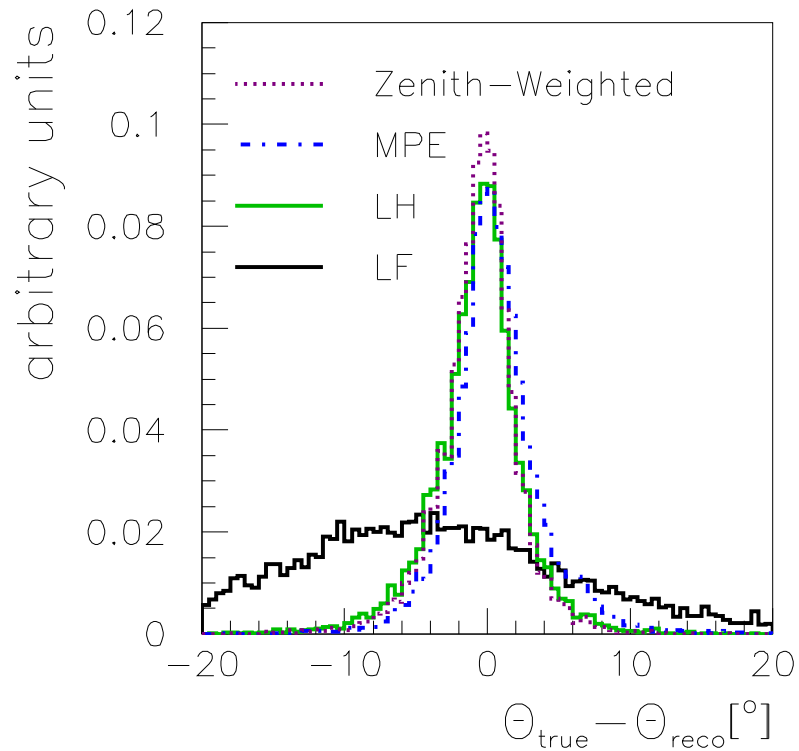
$$\mathcal{L}(E \equiv \{t_1, \dots, t_n; \rho_1, \dots, \rho_n\} \mid H \equiv \{x, y, z, \theta, \phi\})$$

is given by the track geometry and the time residual function (tabulated photon simulation or Pandel function)

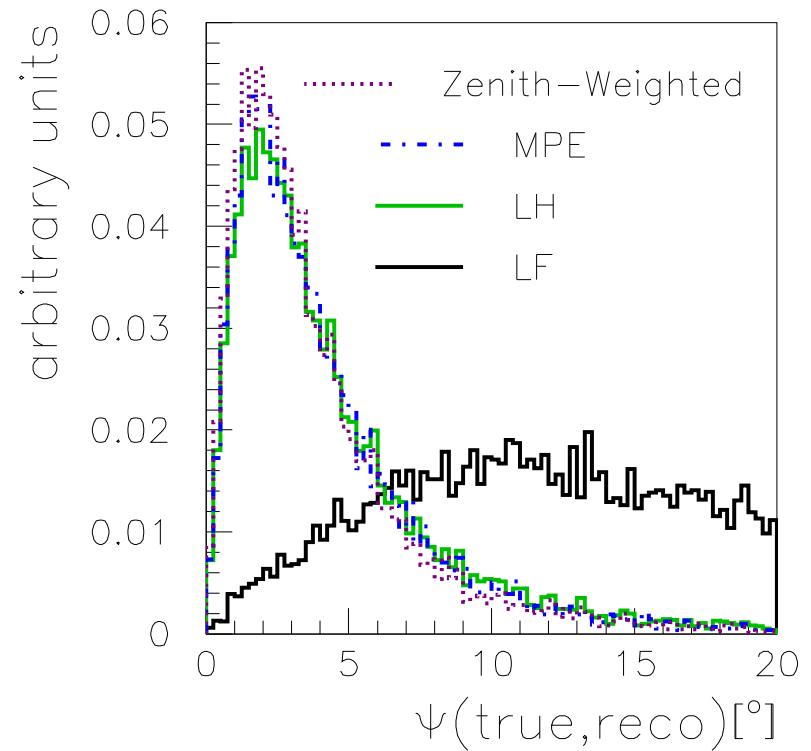
- use a minimisation algorithm (Nelder-Mead simplex, Powell's gradient descent, Minuit) to fit the track parameters



# Reconstruction performance



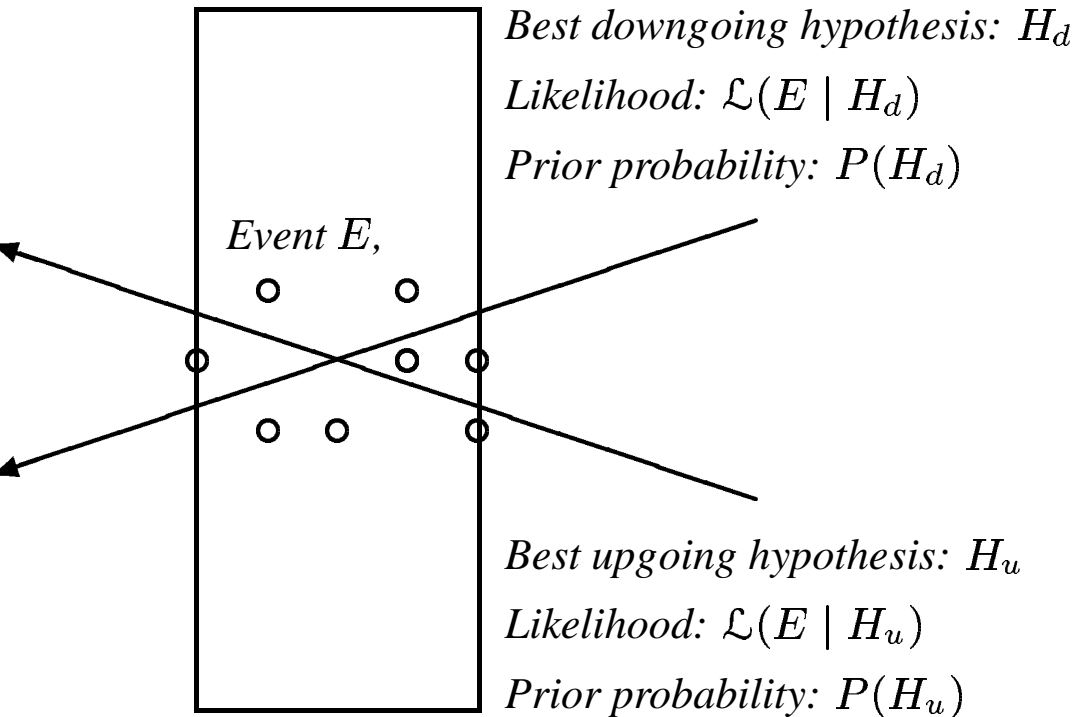
zenith angle



space angle

- Pandel function and photon tables yield similar results

# What is the most probable origin of an observed event?



What if  $\mathcal{L}(E | H_u)$  is only slightly better than  $\mathcal{L}(E | H_d)$ ?

- Should we still choose  $H_u$  over  $H_d$ ?
- We know that  $P(H_d) > P(H_u)$  i.e. more downgoing muons passing through the detector
- Also strong zenith dependence of  $P(H_d)$
- how is this accounted for?

# Joint conditional probability distribution

Most probable downgoing muon hypothesis is the one that maximises joint probability distribution

$$P(E | H_d) P(H_d) \equiv \mathcal{L}(E | \mu_d) \Phi(\mu_d)$$

where  $P(H_d) \equiv \Phi(\mu_d)$ , the flux of downgoing muons in the vicinity of the detector.

Most probable upgoing hypothesis : maximise

$$P(E | H_u) P(H_u) \equiv \mathcal{L}(E | \mu_u) \Phi(\mu_u)$$

where  $P(H_u) \equiv \Phi(\mu_u)$ , the flux of upgoing muons in the vicinity of the detector (taken as uniform).

# *Zenith weighted reconstruction in practice*

- Treat the downgoing muon prior as a simple function of the zenith angle (polynomial fit to simulated muon flux at the detector)
- For each event, find the maximised downgoing and upgoing likelihoods, then take the ratio.
- Use this ratio as a cut parameter, optimised on simulated downgoing and upgoing events
- Rejection of mis-reconstructed atmospheric muons improved by a couple of orders of magnitude over conventional “all hypotheses are equal” method
- Cuts are simplified (in principle, this is the only cut we need)

## ***Bayesian statistics interpretation***

The probabilities of observing an event  $E$  due to up and downward muons are found by integration over the likelihood and priors

$$P_d(E) = \int_{H_d} \mathcal{L}(E | H_d) P(H_d) dH_d$$

$$P_u(E) = \int_{H_u} \mathcal{L}(E | H_u) P(H_u) dH_u$$

The ratio  $P_d(E)/P_u(E)$  is known as the Bayes' discriminant and is the statistically most powerful separator of classes of hypotheses

## *Are we evaluating the discriminant?*

We have approximated the Bayes' discriminant ratio of integrals by the ratio of the maximum values of the integrands :

$$\frac{\int_{H_d} \mathcal{L}(E | H_d) P(H_d) dH_d}{\int_{H_u} \mathcal{L}(E | H_u) P(H_u) dH_u} \simeq \frac{\mathcal{L}(E | \hat{H}_d) P(\hat{H}_d)}{\mathcal{L}(E | \hat{H}_u) P(\hat{H}_u)}$$

# *Don't most physicists reject Bayesian inference?*

- Absolutely yes when used incorrectly!
- Classic example is in upper limit calculations where uniform priors are used to represent subjective “degree-of-belief” about an unknown physical quantity (e.g. the rate of a Poisson process  $\lambda$ , or the mass of a particle  $m$ )
- After measuring  $x$ , an inference on  $m$  is made from  $P(m | x) \propto P(x | m)P(m)$
- Usually take  $P(m)$  to be uniform in some interval
- However  $P(m)$  uniform does not yield same inference as taking  $P(m^2)$  uniform and both choices of “metric” ( $m$  or  $m^2$ ) are equally valid

# *What about our “Bayesian” reconstruction?*

- Acid test - Advanced Statistics in Particle Physics Workshop, Durham, 2002, Ty DeYoung with an audience of the staunchest Bayesians and anti-Bayesians
- Bayesians naturally said the technique was fine....



## ***Bob Cousins for the frequentists....***

- Bayes theorem applies to all types of probability - both subjective degree of belief (e.g. "I think the mass of the Higgs is uniform in the interval 80-200 GeV") and to classical relative frequency probabilities ("the distribution of cosmic rays arriving at earth is uniform and follows a power law energy spectrum")
- Our muon flux “prior” is a relative frequency probability - it’s very easy to define  $P(\mu_d)$  - the muon flux is well measured, theoretically calculated and understood - not subjective at all
- More explicitly : the procedure is “Bayesian” only in that Bayes’ theorem was used – definitely not “subjective” Bayesian!

# *NEVOD experiment developed this technique independently*

- When presenting this work in 2001 in Hamburg, during discussion time A.A. Petruhkin from the NEVOD experiment explained how they did exactly the same thing...
- ... and where able to separate an atmospheric neutrino candidate from a  $10^{10}$  to one background of atmospheric muons in a tiny ( $6 \times 6 \times 7.5\text{m}^3$ ) surface detector!

# *Modern machine learning classification*

- Machine learning - feed a routine a bunch of labelled signal and background, build a model of the Bayes' posterior for future classification of new data
- Neural networks are an example
- Modern methods - Support Vector Machines, Penalised Likelihood methods (Reproducing Kernel Hilbert Space methods)

## *Penalised likelihood method*

- Build a model of the Bayes' discriminant using weighted sums of basis functions and regularisation methods to control the smoothness of the solution
- Currently building a model of atmospheric and isotropic  $E^{-2}$  neutrinos for our diffuse limit analysis (work in collaboration with UW Statistics)

# Conclusions

- Reconstruction of muon tracks in a scattering medium has been successful
- Methods of optimally classifying events as signal and background have been implemented (zenith weighted reconstruction) and are under development (Penalised Likelihood Estimate model building)
- These provide a unifying framework for the reconstruction and classification problem